

Perspective



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# Self-trapped Excitons in All-inorganic Halide Perovskites: Fundamentals, Status and Potential Application

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ABSTRACT: Photoluminescence is a radiative recombination process of electron hole pairs. Self-trapped exciton (STE), occurring in a material with soft lattice and strong electron-phonon coupling, emits photons with broad spectrum and

large stokes-shift. Recently, series halide perovskites with efficient STE emission have been reported and showed promise for solid state lighting. In this Perspective, we present an overview of various photoluminescence with the emphasis on the mechanism and characteristics of emission derived from STE. This is followed by the introduction of STE emission in hybrid halide perovskites. We then introduce all-inorganic STE emitters and focus in particular on the mechanism of STE in double perovskite Cs<sub>2</sub>AgInCl<sub>6</sub> and strategies for efficiency improvement. Last we summarize the current photoluminescence and electroluminescence applications of STE emitters as well as the potential in luminescent solar concentrators, and overlook future research opportunities.

#### **TOC GRAPHICS**



**KEYWORDS**: self-trapped exciton, lead-free perovskite, low dimension, large stokes-shift, broad spectrum.

Photoluminescence (PL) is a basic and important property of semiconductors. There are two types of electronic transitions accompanied with photoluminescence in solids: intrinsic photoluminescence and extrinsic photoluminescence.<sup>1</sup> Intrinsic photoluminescence can be further divided into band-to-band luminescence, exciton luminescence and cross luminescence. Extrinsic luminescence further comprises unlocalized type, such as donoracceptor pair luminescence as well as luminescence caused by isoelectronic trap, and localized type like ionoluminescence and defect luminescence.<sup>1</sup>



**Figure 1.** (a) Schematic of various intrinsic photoluminescence, including bandto-band luminescence, exciton luminescence, STE luminescence as well as cross luminescence. (b) Schematic of the energy level structure of STE (GS = ground state, FE = free exciton state, FC = free carrier state, STE = selftrapped exciton state,  $E_g$  = bandgap energy,  $E_b$  = exciton binding energy,  $E_{st}$  = self-trapping energy,  $E_d$  = lattice deformation energy,  $E_{PL}$  = emission energy). (c) Schematic of the nonradiative recombination process for STE when *S* is

large. Orange circles represent excited electrons.

As shown in Figure 1a, band-to-band luminescence indicates that electrons in the conduction band recombine with holes in the valence band. The emission of mainstream light-emitting-diodes belongs to band-to-band luminescence<sup>2</sup>. Cross luminescence will appear when an electron in the valence band and a hole in the core band recombine. This kind of luminescence is often observed in alkali halides and alkaline-earth halides, and owing to the fast recombination time (even <1 ns)<sup>3</sup>, it plays a crucial role in fast scintillators. Exciton luminescence means that an electron and a hole attract each other by Coulomb interaction and then recombine, emitting a photon. For semiconductors with low exciton binding energy, exciton luminescence can only be observed at low temperature as excitons will convert to free carriers with the aid of thermal energy at room temperature. The exciton emission energy is normally slightly lower than the bandgap by the exciton binding energy, but there is an unusual kind of exciton, called self-trapped exciton (STE), whose emission energy is much smaller than the bandgap<sup>1, 4</sup>.

STEs widely exist in halide crystals<sup>5-6</sup>, condensed rare gases<sup>7</sup>, organic molecular crystals<sup>8</sup>. In these solids, electron-phonon interactions are strong enough for excited electrons and holes to cause elastic distortions in the lattice surrounding them. Thus once electrons and holes are photo-generated, they will be quickly self-trapped because the self-trapped state is more stable than the one in which they would move, dragging the lattice distortion. For example, in alkali halides photo-generated holes are self-trapped by two adjacent halogen ions approaching one another. This molecular ion type center  $X_2^-$  is called a V<sub>k</sub> center and eventually emits intrinsic luminescence when further captures an electron<sup>5</sup>. In the process of STE, exciton will lose some energy, called selftrapping energy  $E_{st}$ . At the same time the energy of ground state will rise due to lattice deformation, and the increased energy is called lattice deformation energy  $E_{d}$ . Consequently, the emission energy can be described by  $E_{PL}=E_{q}-E_{b}$ - $E_{sr}E_d$  (Figure 1b), which is the origin of the large stokes-shift in STE emission.

Extrinsic luminescence is also very common for many practical photoluminescence applications. The mechanism of most phosphors belongs to

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extrinsic luminescence. For instance, green phosphor ZnS with Cu,Al doping is a typical donor-acceptor pair luminescence material<sup>9</sup>. Green light-emitting diode GaP with N isoelectronic doping and red light-emitting diode GaP with Zn-O isoelectronic doping are representative isoelectronic trap luminescence, which were widely applied to the market before the breakthrough of InGaN and AlGaInP light-emitting diodes<sup>2</sup>. Ionoluminescence is also one important sort of extrinsic luminescence, as exhibited in rare earth ions. Rare earth ions are famous for their efficient and stable luminescence with tunable and generally broad spectrum depending on the host lattice, which is necessary for high quality phosphor used in general lighting.

However, in view of the scarcity of rare earth resources and the importance of rare earth element to high-performance magnets, catalysts, alloys, glasses, and electronics, it is highly significant to develop high performance rare-earth free phosphors. In particular, single-source phosphors with white emission not only avoid the self-absorption and color instability problems associated with the current method of combining multiple phosphors for white

emission, but also simplifies the fabrication procedure Broadband and white emission (full width at half maximum, FWHM>100 nm, covering the entire visible spectrum) emitters are extremely rare for most intrinsic luminescence. Encouragingly, STE emission generally features broad spectrum. As mentioned before, strong electron-phonon coupling is essential for the STE formation, which is evaluated by the well-known Huang-Rhys factor S.<sup>4</sup> The electronphonon coupling has a strong connection with the FWHM of luminescence, as described by the equation<sup>10</sup>:

$$FWHM=2.36\sqrt{S}\hbar\omega_{phonon}\sqrt{coth\frac{\hbar\omega_{phonon}}{2k_bT}}$$
 (1)

Where  $\hbar$  is reduced Planck constant,  $\omega_{phonon}$  is phonon frequency, *T* is temperature, and  $k_p$  is Boltzmann constant, respectively. In principle, *S* can be obtained by fitting the temperature dependent *FWHM* of PL peaks with the equation above. Although there is no absolute functional relation between the existence of STE and the value of *S*, it is certain that the larger *S* value is, the more likely STE forms. On the other hand, the coordinate difference

between the free-exciton and STE configurations,  $\Delta Q$  has positive correlation

with S.<sup>4</sup> An overly large S can bring us a broad emission but also lead to the tendency for the excited- and ground-state curves to cross in the configuration coordinate diagram, which means that some excited electrons and holes can recombine via this interaction non-radiatively, emitting several phonons (Figure 1c).Thereby, S should be a suitable value, neither too big nor too small to achieve an efficient STE emission.

Since 2014, a number of halide perovskites have been found to emit abnormally broad spectra with a strength visible to naked eyes.<sup>11-20</sup> Through the replacement of A-site and X-site elements the spectrum can be continuously adjusted from cold-white to warm-white. This kind of perovskites with great potential lighting applications have aroused wide attention and further researches indicated that STE is responsible for the unique emission.<sup>21-23</sup> It is worth noting that several all-inorganic perovskites with excellent STE emitting were reported recently.<sup>24-26</sup> In this Perspective, we will overview the basic property of organic halide perovskites with STE emission. We will then highlight

the milestones achieved in the optical property and the mechanism of the leadfree all-inorganic perovskites with efficient STE emission. Finally, we will discuss the photoluminescence and electroluminescence (EL) based applications of all-inorganic lead-free perovskites with STE emission as well as future research opportunities.



**Figure 2.** (a) The photograph of  $(EDBE)[PbBr_4]$  under 365nm light, inset shows the perovskite under ambient light (EDBE = 2,2)-

(ethylenedioxy)bis(ethylammonium)). Reprinted with permission from ref 15.

Copyright 2014 American Chemical Society. (b) X-ray structure of EDBE[PbBr<sub>4</sub>]:

gray, green, red, brown, blue represent Pb, Br, O, C and N atoms. H atoms

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are omitted for clarity. (c) Absorption (black) and emission (blue) for  $(AEA)PbBr_4$  (AEA = 3-(2-ammonioethyl)anilinium). (d)Schematic of interoctahedral tilting  $\theta_{tilt}$  in a single Pb-Br layer. (e)Values of  $\ln(I_{BE}/I_{NE})$  for various layered lead halide perovskites as a function of interoctahedral tilting  $\theta_{tilt}$ .  $D_{tilt}$  represents 180° -  $\theta_{tilt}$ ; BE and NE represent broad emission and narrow emission in layered perovskites. Panels c-e are reprinted with permission from ref 22. Copyright 2017 American Chemical Society. (f) The photograph of 0D tin and antimony perovskites under ambient light and 365nm light. (g) Crystal structure of (C<sub>4</sub>N<sub>2</sub>H<sub>14</sub>Br)<sub>4</sub>SnBr<sub>6</sub>: red, green, blue, gray represent Sn, Br, N, C atoms; H atoms are hidden for clarity. (h) PL and PLE of  $(C_4N_2H_{14}Br)_4SnBr_6$ , (C<sub>4</sub>N<sub>2</sub>H<sub>14</sub>I)<sub>4</sub>SnI<sub>6</sub>, (C<sub>9</sub>NH<sub>20</sub>)<sub>2</sub>SbCl<sub>5</sub>. Panels f-h are reprinted with permission from ref 19 published by The Royal Society of Chemistry.

Karunadasa et al. and Kanatzidis et al. first reported several low dimensional lead halide perovskites with intrinsic broadband emission from 400

nm to 800 nm, covering the visible region, such as (N-MEDA)[PbBr<sub>4-x</sub>Cl<sub>x</sub>] (N-MEDA = N-methylethane-1,2-diammonium, x =  $0^{-1.2}$ , EA<sub>4</sub>Pb<sub>3</sub>Cl<sub>10-x</sub>Br<sub>x</sub> (EA = ethylammonium, x<10) (Figure 2a).<sup>14-16, 18, 21-23</sup> All of them include a large organic group at A-site and further form a two-dimensional (2D) structure (Figure 2b). Thus compared with three-dimensional (3D) perovskites they all have an obvious exciton absorption peak (Figure 2c). Please note that in this Perspective, all the perovskite dimension refers to the connection of octahedron instead of material size. Due to the large organic groups and the mixed halides in some cases, the lead octahedra and the inorganic frameworks show distortions to some degree, which can be describe by the interoctahedral tilting  $\theta_{tilt}$  (Figure 2d).<sup>14, 18</sup> The STE property depends profoundly on the inorganic structure distortion. Besides, further research showed that there exists a potential barrier for self-trapping in two dimensional systems, which could allow some excited carriers to recombine as free exciton (FE) luminescence rather than STE luminescence.<sup>21</sup> The emission spectra of 2D perovskites usually comprise two peaks (Figure 2c): the main broad emission with lower energy is

STE luminescence and the secondary narrow emission with higher energy is

FE luminescence. The ratio of FE luminescence and STE luminescence is also dependent on the distortion (Figure 2e).<sup>23</sup> A warm-white STE luminescence and a blue FE luminescence constitute a standard white emission. Thus, white-light emitter can be designed with appropriate A-site organic group<sup>23</sup>. But as mentioned before, large S, and probably some deep defects in those perovskites results in a low emission efficiency. The highest PLQY in lead halide perovskites with white emission is 9% in (EDBE)PbBr<sub>4</sub>.<sup>23</sup> For the detailed discussion on STE in this type of materials, readers are encouraged to refer to the excellent review written by Prof. Karunadasa.<sup>23</sup>

In parallel, Ma et al. and Ning et al. reported a series of lead-free tin and antimony perovskites with broadband emission and especially a near unity photoluminescence quantum yield (PLQY), which have outstanding efficiency and avoid the environmental issues caused by Pb (Figure 2f).<sup>13, 17, 19-20, 27</sup> For those perovskites, the B-site element is small while the A-site component is large in volume. Therefore the  $[SnBr_6]^{4-}$  or  $[SbCl_5]^{2-}$  is totally isolated by the

large organic group, forming so-called 0D structure (Figure 2g). Photogenerated carriers are confined in these isolated octahedral, experiencing strong quantum confinement and hence enjoying high PLQY. In 0D systems, there is no interconnected inorganic framework and hence no conception of interoctahedral tilting which plays a significant role in the emission of 2D lead halide perovskites.<sup>22</sup> Thus the STE property depends mainly on the distortion in a single octahedron and the octahedral composition. The energy barrier between FE and STE is believed not to exist in 0D materials.<sup>28-30</sup> This is probably because the immediate distortion of the more mobile [SnBr<sub>6</sub>]<sup>4-</sup> or [SbCl<sub>5</sub>]<sup>2-</sup> octahedra, leading to sole STE emission. Therefore contrast to 2D lead halide perovskites with STE emission discussed above, these 0D lead-free perovskite have high emission efficiency but a narrower emission with very good Gaussian shape (Figure 2h) because of the absence of FE emission.<sup>19</sup> The disappearance of FE luminescence means that 0D perovskites are capable to own warm-white emission at most unless the FWHM of STE emission is larger than 250 nm, which is at the expense of PLQY and unworthy.

STE formation is believed to have a strong relationship with the dimensionality of the system.<sup>21, 28-32</sup> In other words, a large electron-phonon coupling is crucial to the formation of STE, and the low crystal dimensionality can facilitate octahedral distortion and STE formation, as well as high emission efficiency because of the strong quantum confinement particularly in the 0D perovskites. Despite diverse organo-ammonium cations are available for A site, their moisture and thermal instability restricts practical applications. Thereby, all-inorganic perovskites are proposed as a promising alternative considering their excellent stability.



**Figure 3.** (a) Crystal structure of  $Cs_4SnBr_6$ , grey, green and violet represent Sn, Br and Cs atoms. (b) The PL and PLE of  $Cs_4SnBr_6$ . (c) The photograph of  $Cs_4$ .  ${}_xA_xSn(Br_{1-y},I_y)_6$  (A=K, Rb) under 365nm light. Panels a-c are reprinted from ref 24 with permission from John Wiley & Sons Ltd. (d) The photograph of  $Cs_3Cu_2I_5$  crystal under ambient light, inset shows the crystal under 245nm light. (e) Crystal structure of  $Cs_3Cu_2I_5$ , blue, violet and cyan represent Cu, I, Cs atoms. (f) The photoluminescence (PL) and photoluminescence excitation (PLE) of  $Cs_3Cu_2I_5$ . Panels d-f are reprinted from ref 25 with permission from John Wiley & Sons Ltd.

In 3D all-inorganic perovskite with crystal structure of cubic phase ABX<sub>3</sub>, crystal lattice is difficult to distort because octahedra connect with each other via vertex firmly. Even formed, the STE is unstable.<sup>33</sup> As discussed before, a reduced dimensionality and localized electrons and holes contribute to the formation of efficient STE. For example, yellow orthorhombic phase CsPbl<sub>3</sub> is a kind of low dimensional inorganic perovskite and thus STEs emission was observed.<sup>34</sup> Recently Kovalenko et al. reported 0D material Cs<sub>4</sub>SnBr<sub>6</sub> with a relatively high PLQY (~15%).<sup>24</sup> From the perspective of crystal structure, [SnBr<sub>6</sub>]<sup>4-</sup> octahedron is separated by Cs<sup>+</sup> ions totally (Figure 3a). The optical property is also a typical STE emission: large stokes-shift and broad spectrum (Figure 3b). Simulation showed Jahn-Teller distortion in [SnBr<sub>6</sub>]<sup>4-</sup> is the origin of STE.<sup>24</sup> Compared with those hybrid tin 0D perovskites discussed above, the emission spectrum is slightly broader but PLQY is obviously lower. However, the moisture and thermal stability of Cs<sub>4</sub>SnBr<sub>6</sub> is significantly enhanced because of the absence of organic group. Another superiority of  $Cs_4SnBr_6$  is the optical

property can be adjusted continuously by substituting A site with K<sup>+</sup> or Rb<sup>+</sup> and X site with I<sup>-</sup> (Figure 3c).<sup>24</sup> But the frangibility to oxidation of Sn<sup>2+</sup> remains a significant concern for long-term stability. Thus it is meaningful to find more stable and efficient all-inorganic STE emitter.

Hosono et al. reported a nontoxic all-inorganic halide 0D material Cs<sub>3</sub>Cu<sub>2</sub>I<sub>5</sub> with efficient broad blue emission (PLQY up to 90%) (Figure 3d).<sup>25</sup> Similarly, the emission center,  $[Cu_2l_5]^{3-}$ , is isolated with each other by Cs<sup>+</sup> in  $Cs_3Cu_2I_5$  (Figure 3e). Thus excitons are spatially confined around emission centers, favoring the radiative recombination and contributing a high PLQY. In addition, the efficient blue emission spectrum of  $Cs_3Cu_2l_5$  is comparatively narrow among perovskites with STE emission (Figure 3f), which is also rare and of great value. Furthermore, Cu<sup>+</sup> is thermodynamically stable and more anti-oxidative contrast to Sn<sup>2+</sup>; however its moisture stability could still be a concern. The ideal STE emitter should combine a near-unity PLQY with an excellent air, light and thermal stability, and preferably white emission as for solid state lighting applications. Double perovskites emerged as the potential

substitute for lead halide perovskites.<sup>35-44</sup> The elements making up double perovskites are Ag<sup>+</sup>, Bi<sup>3+</sup> and In<sup>3+</sup> etc.<sup>37, 40, 42</sup> Those elements generally have fixed chemical valence, ensuring better stability. Moreover, recent study indicated that despite the 3D structural dimensionality of double perovskites, many of them have a 0 D electronic dimensionality, which is a shortcoming for solar cell absorbers but a merit for STE emitter.<sup>35, 45</sup>





from ref <sup>46</sup>. Copyright 2017 American Chemical Society. (b) STE in Cs<sub>2</sub>AgInCl<sub>6</sub>.

The cyan and magenta isosurfaces represent the electron states and the hole states. Inset shows the Jahn-Teller distortion of the [AgCl<sub>6</sub>] octahedron. (c) The parity change of electron wave function (isosurface at Ag atom) of STE between  $Cs_2AgInCl_6$  and  $Cs_2Na_xAg_{1-x}InCl_6$ . (d) The confined STE in  $Cs_2Na_xAg_{1-x}InCl_6$ . (e) The absorption and photoluminescence of  $Cs_2AgInCl_6$ ,  $Cs_2NaInCl_6$  and  $Cs_2Na_0.4Ag_{0.6}InCl_6$ . (f) Transient absorption spectra for  $Cs_2Ag_{0.6}Na_{0.4}InCl_6$ . (a) Emission intensity versus excitation power for  $Cs_2Ag_{0.6}Na_{0.4}InCl_6$ . Panels b-g are reprinted from ref 26 with permission from Springer Nature.

In our previous work we found double perovskite Cs<sub>2</sub>AgInCl<sub>6</sub> exhibits a warm-white emission spectrum from 400 nm to 800 nm.<sup>46</sup> Besides the main broad peak at 590 nm, a weak peak shows up at 430 nm (Figure 4a), which may be attributed to FE emission. The lifetimes of two peaks are also in agreement with FE and STE emission.<sup>46</sup> Simultaneous emission of FE and STE is the feature of 2D and 3D perovskites with STE emission, which has been discussed before. Theoretical calculation revealed that due to the delocalized In

5s states, the wave function of electrons is dispersive and has a 3D structure, while the holes are confined at [AgCl<sub>6</sub>] with a 0D structure (Figure 4b).<sup>26</sup> After excitation, holes will be partially trapped at Ag atoms quickly and change the electronic configuration from 4d<sup>10</sup> to 4d<sup>9</sup>, favoring a Jahn-Teller distortion and further resulting in the STE formation as the process in the formation of STE in AgCI.<sup>26</sup> FE emission still present in Cs<sub>2</sub>AgInCl<sub>6</sub>, in analogy with 2D or 3D hybrid halide perovskites, enabling potential white emission. We obtained the Huang-Rhys factor in Cs<sub>2</sub>AgInCl<sub>6</sub> with theoretical value 37 and experimental value 38.7.<sup>26</sup> This is an appropriate value for efficient emission but the PLQY of  $Cs_2AgInCl_6$  is extremely low (<0.1%). One reason is that the FE and STE are parity-forbidden transition. Another reason is small overlap between the wave function distributions of electrons and holes (Figure 4b).<sup>26</sup> All in all, Cs<sub>2</sub>AgInCl<sub>6</sub> is a promising white emission material as long as these two problems are overcome.

Recently we made a breakthrough by Na alloying Cs<sub>2</sub>AgInCl<sub>6</sub>. Parityforbidden transition derives from the inversion symmetry of the crystal structure

in Cs<sub>2</sub>AgInCl<sub>6</sub>.<sup>47</sup> Thereby we proposed that the parity-forbidden property could be removed as long as the inversion symmetry was broken, for example, via partial substitution to Ag. Alkali mental has a distinctively different electronic configuration to Aq and thus is an appropriate candidate. In addition,  $Cs_2AgInCl_6$  and  $Cs_2NaInCl_6$  have the same crystal structure and a negligible lattice mismatch (0.3%).<sup>26</sup> We thus anticipated, and experimentally validated, that there would not be detrimental defects or phase separation when Na was introduced into  $Cs_2AqInCl_6$ . Theoretical calculation revealed that in  $Cs_2Na_xAq_{1-1}$ <sub>x</sub>InCl<sub>6</sub>, the STE is spatially confined by Na<sup>+</sup>, leading to enhanced electron and hole orbital overlap (Figure 4d).<sup>26</sup>The inversion symmetry of the Cs<sub>2</sub>AgInCl<sub>6</sub> is also broken upon Na alloying (Figure 4c). These two factors combined resulted in the improvement of luminescence intensity by 3 order of magnitude. Experimental evidence for the decreased electronic dimensionality is that contrast to the origin Cs<sub>2</sub>AgInCl<sub>6</sub>, Na alloyed Cs<sub>2</sub>AgInCl<sub>6</sub> exhibits an evident exciton absorption peak at 365nm (Figure 4e). Na alloying also raised the activation energy and decomposition temperature, and thus helping the emission

stability. Trace  $Bi^{3+}$  was further introduced to diminish defects in  $Cs_2Na_xAg_{1-}$ <sub>x</sub>InCl<sub>6</sub> and depress the non-radiative recombination and further enhance PLQY to 86% (peak value, 70% for average value), an unprecedented emission efficiency for white emissive phosphors.

In addition, STE emission is also observed in double perovskites nanocrystals. Manna et al. first reported Cs<sub>2</sub>AgInCl<sub>6</sub> nanocrystals with white emission (1.6% PLQY)<sup>48</sup>, which exhibited good stability under air. Han et al. also reported a bright dual-color emission in Cs<sub>2</sub>AgBi<sub>x</sub>In<sub>1-x</sub>Cl<sub>6</sub> double perovskite nanocrystals<sup>49</sup>, which is similar with the STE emission mechanism in Cs<sub>2</sub>Na<sub>x</sub>Ag<sub>1-x</sub>InCl<sub>6</sub> bulk materials. Moreover, self-trapping process is also observed in some new perovskite nanocrystals that have not been reported for the bulk form, such as Cs<sub>2</sub>AgSb<sub>x</sub>In<sub>1-x</sub>Br<sub>6</sub><sup>50</sup>. Thus STEs in halide perovskites nanocrystals deserve further researches to deepen the understanding and realize its full potential.

The STEs and defects are apparently similar in many ways; however the crucial difference is that the ground state of STE is actually the perfect lattice.<sup>4</sup> After the free carriers

generated, the lattice will quickly distort and STE will form simultaneously. Generally, the exciton self-trapping time should be comparable to the vibrational period of this phonon mode. Hence the time can be evaluated as  $\tau=2\pi/\Omega$  ( $\Omega$  is phonon energy). This estimated time is about several hundred femtoseconds and has been demonstrated experimentally by transient absorption (Figure 4f).<sup>23, 26</sup> Another accessible experimental evidence is the emission intensity has a linear relationship with excitation power (Figure 4g).<sup>26</sup> That is because the defects emission is easier to be saturated under high excitation power. Other discernable evidence for the STE emission include similar photoluminescence excitation spectrum and photoluminescence lifetimes for emitted photons of different wavelength, and photo-introduced absorption in the transient absorption spectrum<sup>26</sup>.



**Figure 5.** (a) The photograph of Cs<sub>3</sub>Cu<sub>2</sub>I<sub>5</sub> and yellow phosphor and their mixture at the ratio of 2:8 under ultraviolet (UV) light. Reprinted from ref 25 with permission from John Wiley & Sons Ltd. (b) The down-conversion light-emitting-diode (LED) based on Cs<sub>2</sub>Na<sub>0.4</sub>Ag<sub>0.6</sub>InCl<sub>6</sub> phosphor. (c) The current-densityluminescence-voltage curves of Cs<sub>3</sub>Cu<sub>2</sub>I<sub>5</sub> LED. The inset shows the LED device while working. Reprinted from ref 25 with permission from John Wiley & Sons Ltd. (d) The current-density-luminescence-voltage curves of Cs<sub>2</sub>Na<sub>0.4</sub>Ag<sub>0.6</sub>InCl<sub>6</sub> LED. The

inset shows the LED device while working. Reprinted from ref 26 with permission from Springer Nature.

For phosphors applied in general lighting, broad spectrum is required to achieve a high color rendering index (CRI), which is an important indicator for light source to reveal the colors of various objects faithfully.<sup>51</sup> And large stokesshift promises a negligible self-absorption, as self-absorption could seriously restrict the total efficiency of the LED.<sup>52</sup> Accordingly, efficient STE emitters are ideal phosphors used in general lighting. White phosphor based on a mixture of Cs<sub>3</sub>Cu<sub>2</sub>I<sub>5</sub> and yellow phosphor was demonstrated and showed a good performance (Figure 5a).<sup>25</sup> We also demonstrated a warm white light emitting diode with one single phosphor based on the optimized Cs<sub>2</sub>Na<sub>x</sub>Ag<sub>1-x</sub>InCl<sub>6</sub> (Figure 5b).<sup>26</sup> Benefiting from the large exciton binding energy and the high light and thermal stability of double perovskite, the phosphor shows no obvious variation for the spectrum or the intensity, when operating at 333 K, or heated at a 150 °C for 1000 hours, or working at 5000 cd/m<sup>2</sup> for 1000 hours.<sup>26</sup> These

positive results probably promise a huge market prospect for high efficient STE emitters.

For lighting applications, the phosphors are generally excited by an electrically pumped light emitting diodes (LEDs). Hence besides a high PLQY, it is necessary for the optimum excitation wavelength to match this range for STE emitter to achieve high power efficiency. Current III-Nitride based UV LEDs shows excellent performance in the range of 365 nm~400 nm, yet substantially lower efficiency below 365 nm.<sup>53</sup> The optimum excitation wavelength of  $Cs_2Na_xAg_{1-x}InCl_6$  is on the boundary of this range<sup>26</sup> (exactly at 365 nm, Figure 4e), but those of  $Cs_3Cu_2l_5$  and  $Cs_4SnBr_6$  are out of this range<sup>24-25</sup> (Figure 3f and 3b). To reduce the overly large stokes shift, one possible solution is to explore other element alloying or substitution to minimize the lattice deformation energy  $E_d$  and self-trapped energy  $E_{st}$ .

It would be scientifically interesting and technologically attractive to directly make electroluminescent (EL) devices using STE material as the

emission layer. Though many lead and tin halide perovskites with efficient STE emission have been discovered, few STE emission based EL study is reported. One possible reason is the large organic group severely impedes the transport of carriers, and the tremendous heat released during working might destroy the materials themselves. However, this problem can be resolved to a great extent via using all-inorganic STE emitter. EL devices based on both of Cs<sub>3</sub>Cu<sub>2</sub>I<sub>5</sub> and  $Cs_2Na_xAg_{1-x}InCl_6$  have been reported<sup>25-26</sup> (Figure 5c and 5d). As expected, the EL devices emit similar spectra with their PL spectra. Nevertheless, both of them have a large threshold voltage and low efficiency. There are two main obstacles: the low carrier mobility because of their low electronic dimensionality and strong coupling to the lattice, and the poor charge injection due to their large bandgap. Overall, both are fundamental challenges and we believe there remains a long way for STE EL device to realize their full potential.



Figure 6. Schematic of luminescent solar concentrator based on STE emitter.

Considering the high PLQY and large stokes-shift, we propose that these STE emitters could be promising for luminescent solar concentrators (LSCs).<sup>54-59</sup> LSCs are a potential approach to decrease the usage of traditional solar cells by concentrating solar irradiation through luminescent materials (Figure 6). There are mainly two kinds of phosphors used in LSCs: fluorescent dyes and quantum dots (QDs).<sup>59</sup> Fluorescent dyes have large stokes-shift but their stability under long time illumination is unsatisfactory. Quantum dots can have a near-unity PLQY and a good light stability; however, large stoke-shift have to be judiciously engineered<sup>57</sup> (for example, giant shells) and their high-cost

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synthesis is also a concern. STE emitters, particularly Cu<sub>3</sub>Cu<sub>2</sub>I<sub>5</sub> and Cs<sub>2</sub>Na<sub>x</sub>Ag<sub>1</sub>. <sub>x</sub>InCl<sub>6</sub>, combine the advantages of high PLQY, good thermal stability, large stokes shift with minimized self-absorption, and low-cost manufacturing. Solar windows using these STE emitters further enjoy high transparency, an important feature for windows. <sup>58</sup> The shortcoming is these materials often only absorb the UV portion of solar spectrum, thus suffering from low efficiency.

In conclusion, STE derives from large electron-phonon coupling and soft lattice and is featured with broad spectrum and large stokes-shift, making it somewhat peculiar compared with other intrinsic luminescence. 2D hybrid perovskites have co-existence of FE and STE and hence high quality white emission but the obvious setback is low PLQY; 0D hybrid perovskites have a near-unity PLQY but lose the FE component and white emission. The thermal instability brought by organic groups could be improved by the all-inorganic perovskite counterparts. Cs<sub>4</sub>SnBr<sub>6</sub>, Cs<sub>3</sub>Cu<sub>2</sub>I<sub>5</sub> and Cs<sub>2</sub>AgInCl<sub>6</sub> are such three prominent examples. Particularly for Cs<sub>2</sub>AgInCl<sub>6</sub>, through Na alloying to break the parity forbidden transition and reduce the electronic dimensionality and

through  $Bi^{3+}$  doping to suppress non-radiative defects, the optimized white emissive  $Cs_2Na_{0.4}Ag_{0.6}InCl_6$  STE emitters combines outstanding PLQY and excellent thermal, air and light stability and shows promise as down conversion phosphors for solid state lighting .

The study of STE emission in halide perovskites or related variants are only at their infancy stage; there remain a large room to be explored. Instead of V<sub>k</sub> center in alkali halides, Jahn-Teller distortion in octahedra favors the formation of STE in halide perovskites. Further understanding of the lattice distortion process, the decay pathway of STE emission, as well as other relevant mechanism are urgent. Furthermore, all-inorganic perovskites with efficient STE emission are scarce contrast to various organic STE emitting perovskites. Synthesizing new all-inorganic materials with efficient STE emission and preferably high quality white emission with >365 nm excitation is of great significance. So far, the seemly promising application for STE emitting perovskites is phosphor used in general lighting. Therefore exploring new applications, for example, luminescent solar concentrators, is very meaningful.

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#### Notes

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